

Multiplicity Fluctuations in Au+Au Collisions at RHIC

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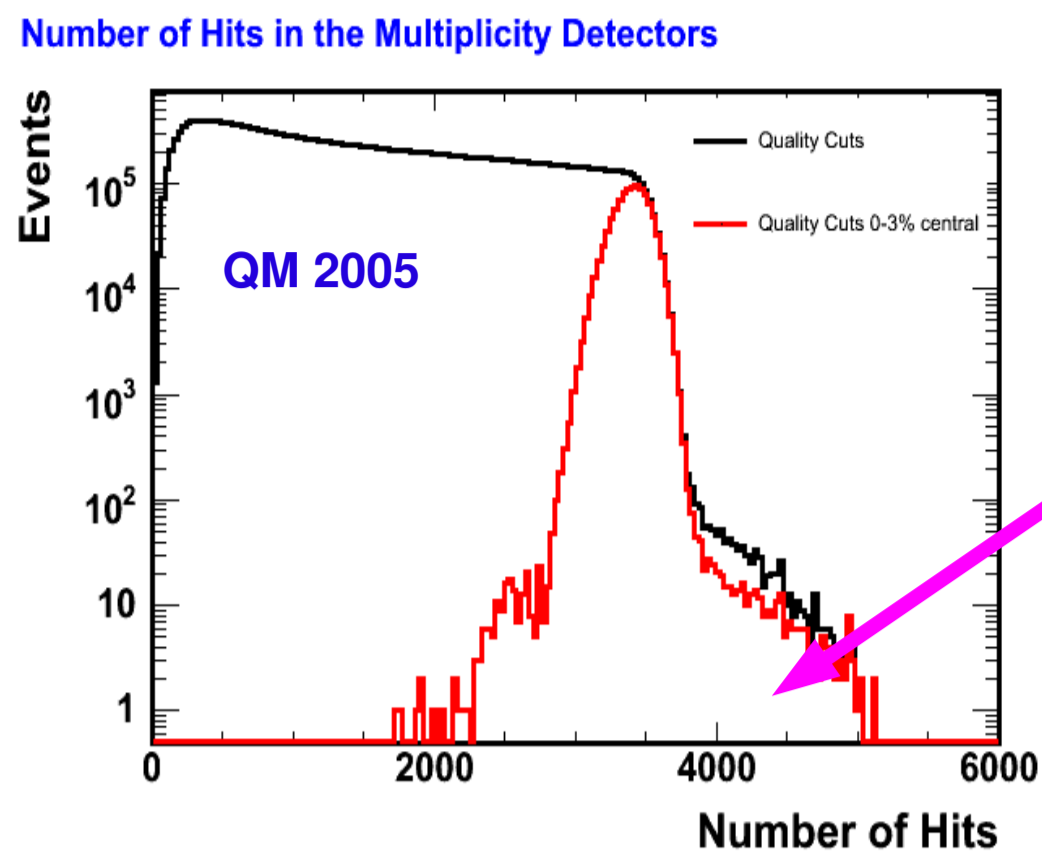
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Search For Enhanced $dN/d\eta$ Fluctuations



Central Au+Au Collisions at $\sqrt{s_{NN}} = 200$ GeV

Tail of large multiplicities consistent with fraction of pileup events
G. Stephans et al., PHOBOS Collaboration QM 2005

Motivation:

Identify unusual events in the most central Au+Au collisions by large deviations of $dN/d\eta$ distribution with respect to normal events

Determination of $dN/d\eta$ for each event:

- silicon pad sensors in the range $|\eta| < 5.3$ used
- number of hits counted in η bins 0.2 wide
- corrections for acceptance and secondary particle production are neglected
- $\langle N(\eta) \rangle$ is calculated in 60x5 bins of Z and Y vertex position

Event by event fluctuations of $dN/d\eta$ are measured by:

$$\chi^2_{NDF} = \frac{1}{N_{bins} - 1} \sum \frac{(N(\eta) - S \langle N(\eta) \rangle)^2}{\sigma^2(N(\eta))}$$

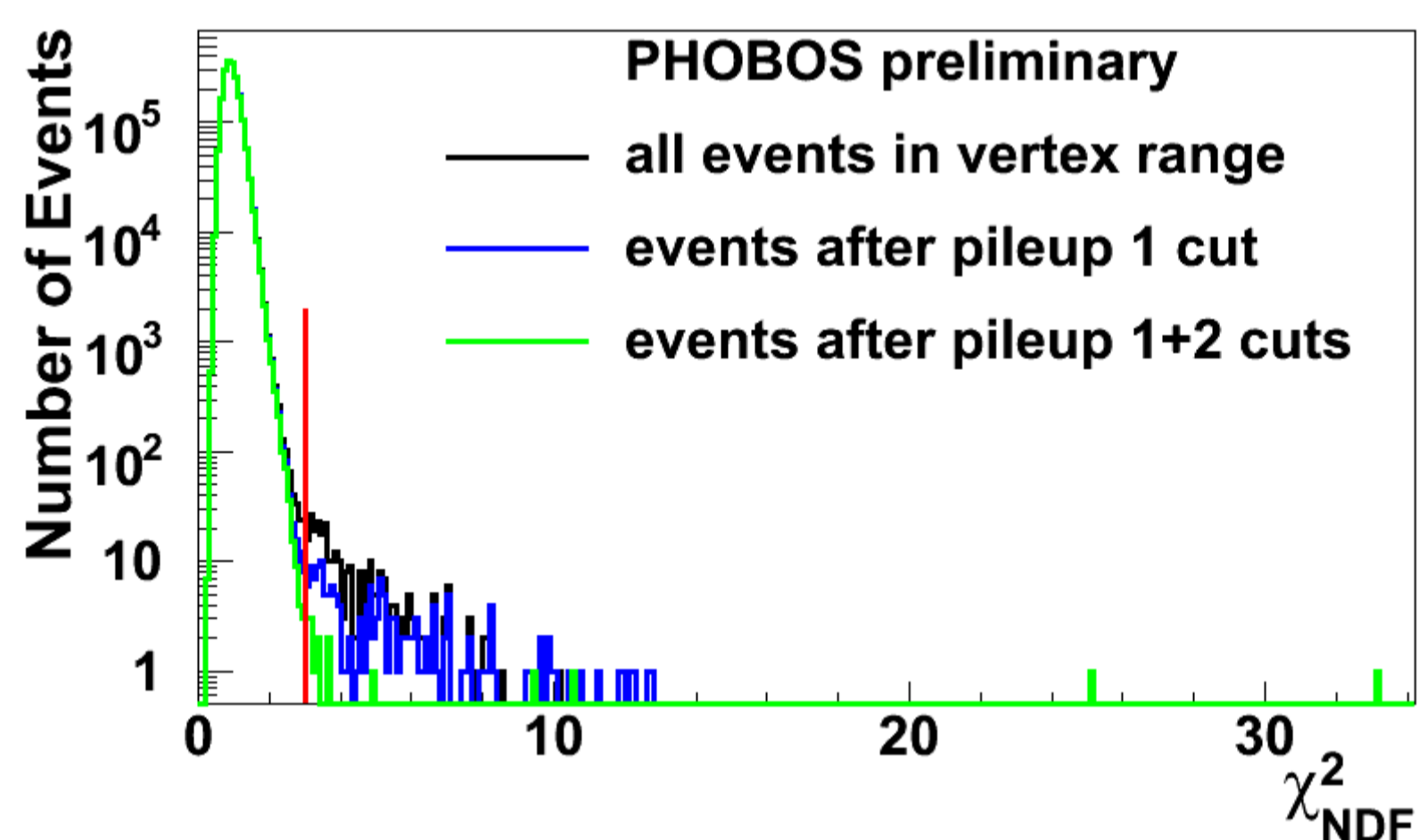
where:

- N_{bins} is the number of eta bins
- $N(\eta)$ is the number of hits in each η bin
- S is the fitted normalization scaling factor, accounting for varying total multiplicity

Events with $\chi^2_{NDF} > 3$ are counted as unusual.

Analyzed sample of central events (about 1,900,000):

- centrality 0-3%, restricted vertex range -10 to 10 cm
- runs or events with hardware problems removed
- two types of pileup removal cuts applied (see below)



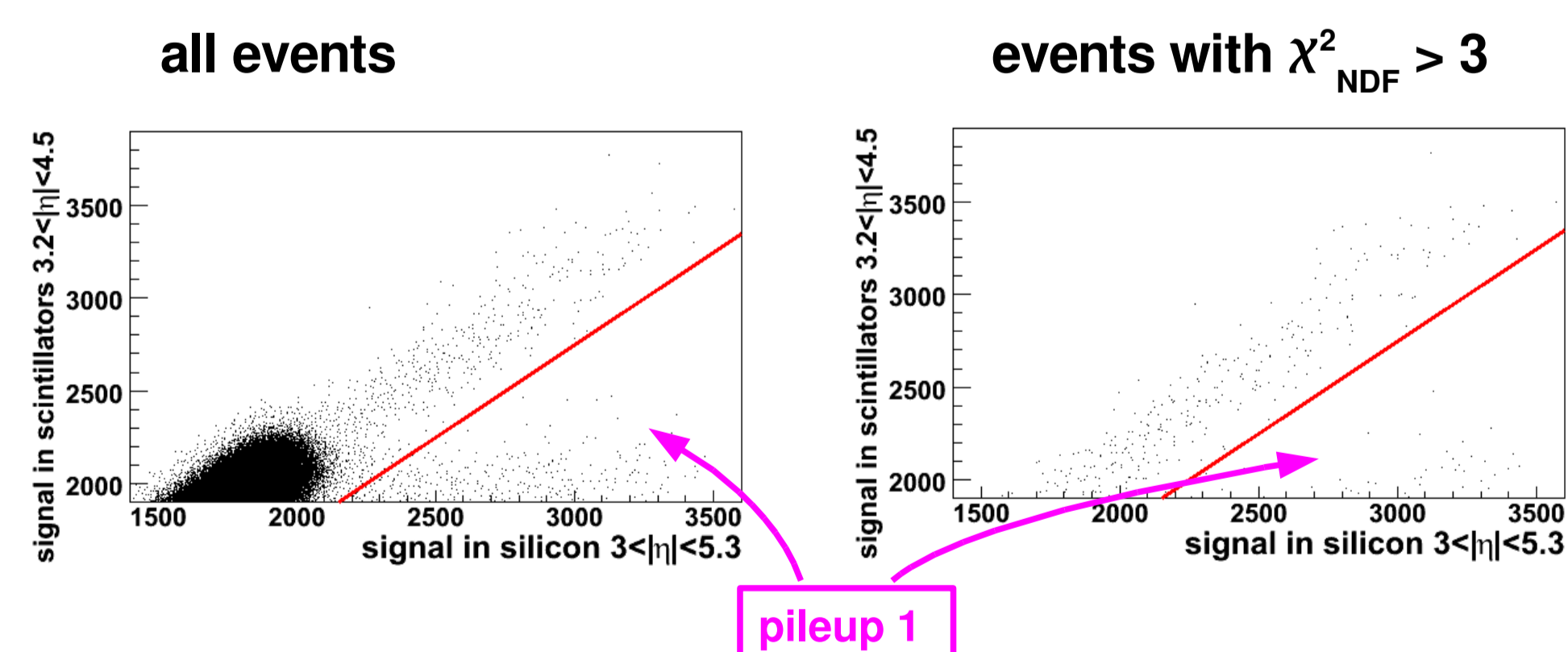
Final result:

The fraction of events with large deviations in $dN/d\eta$ distribution is smaller than 10^{-5}

Identification and removal of pileup events (1)

- two or more events from different bunch crossings:
 - registered as one in silicon sensors due to long integration time
 - registered separately in scintillator counters

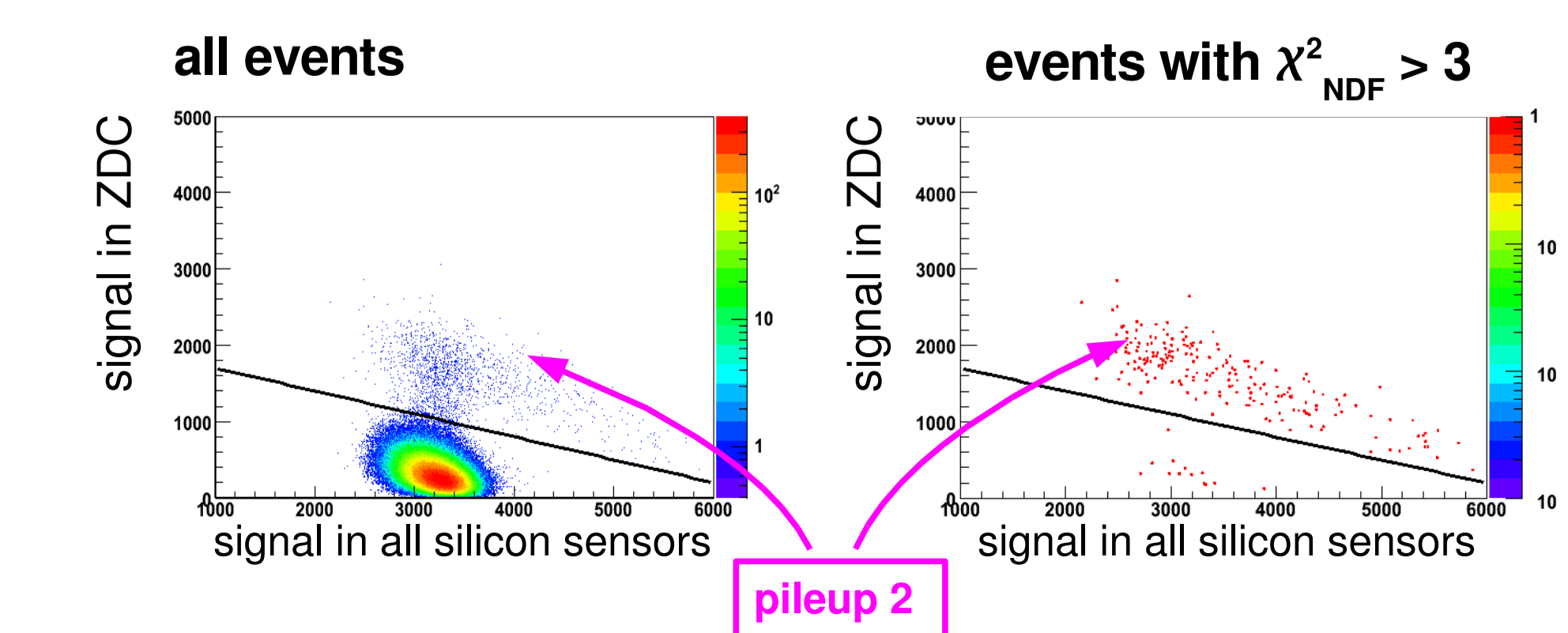
exclude: large signal in silicon detectors and smaller signal in scintillator detector (removes ~50% of events with $\chi^2_{NDF} > 3$)



Identification and removal of pileup events (2)

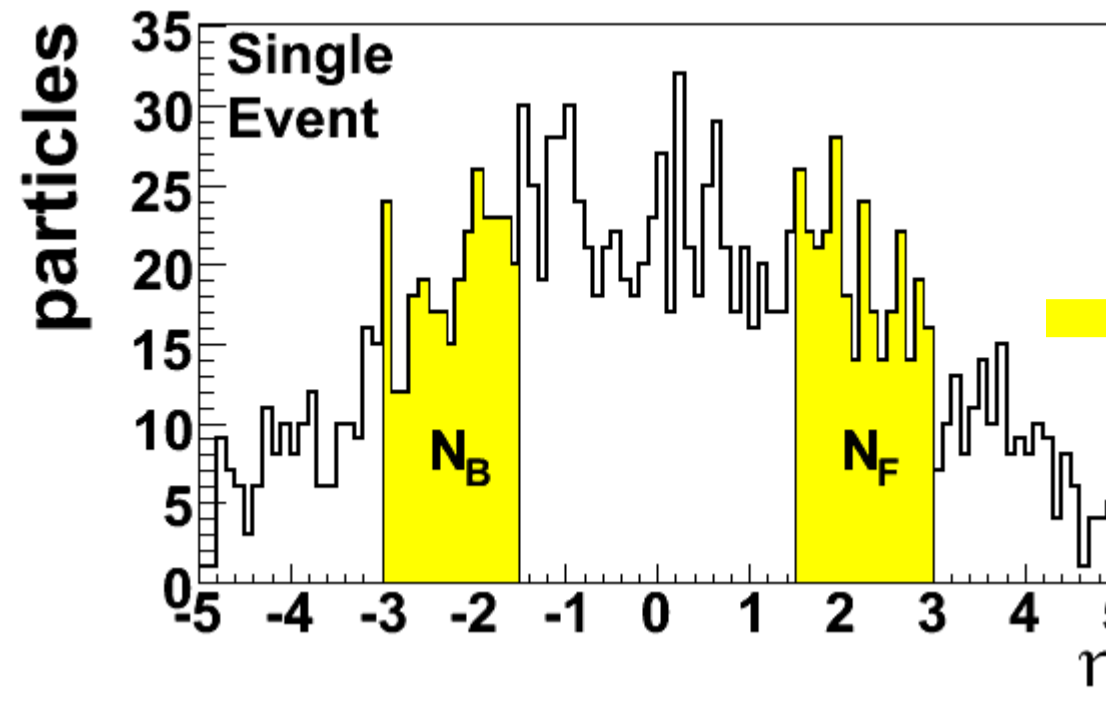
- two events in the same bunch crossing:
 - central event without spectator neutrons (small ZDC signal)
 - less central event with several spectators (larger signal in ZDC)

exclude: large signal both in all silicon sensors ($|\eta| < 5.3$) and in ZDC (removes ~90% of events with $\chi^2_{NDF} > 3$)



Forward-Backward Multiplicity Fluctuations and Cluster Models

Au+Au Collisions at $\sqrt{s_{NN}} = 200$ GeV



$$C = \frac{N_F - N_B}{\sqrt{N_F + N_B}}$$

For statistical fluctuations $\sigma^2(C) = 1$ independent of multiplicity

In Au+Au collisions at highest RHIC energy: $\sigma^2(C) > 1$
Phys. Rev. C74, 011901(R) (2006)
also: K. Woźniak, QM2004, P. Steinberg, QM2005

In the presence of short range correlations caused by production of particles in clusters with multiplicity (size) k , the value of $\sigma^2(C)$ increases:

$$\sigma^2(C) \rightarrow k \sigma^2(C)$$

(for a distribution of k we should use $k_{eff} = \langle k \rangle + \sigma^2(k)/\langle k \rangle$)

Possible explanation:

production of particles in clusters

Two models of cluster production are studied:

Isotropic Cluster Model has one parameter, effective cluster size k_{eff} which can be fitted to the data:

- $k_{eff} = 3.4$ for most central (0-20%) Au+Au collisions
- $k_{eff} = 4.3$ for peripheral (40-60%) Au+Au collisions

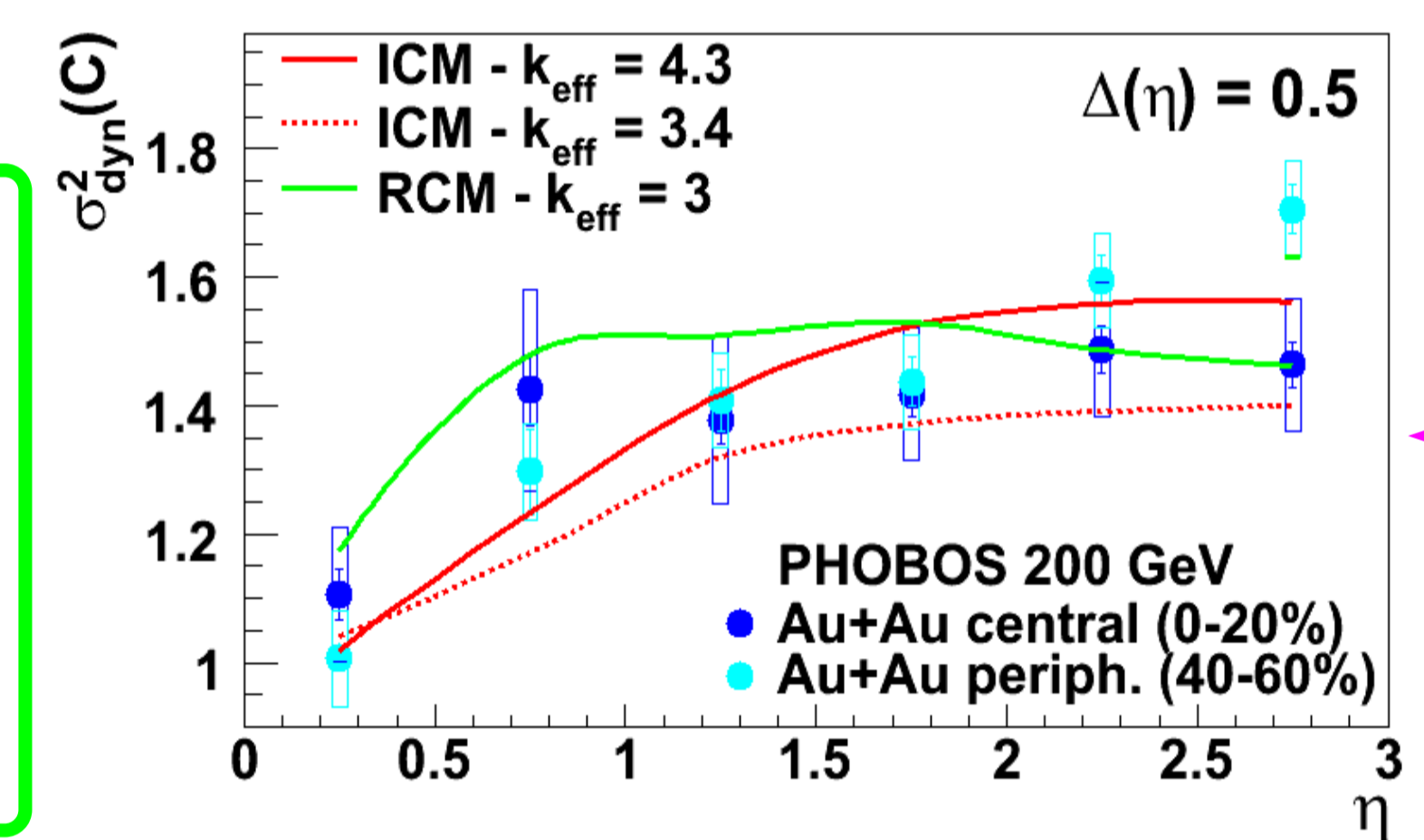
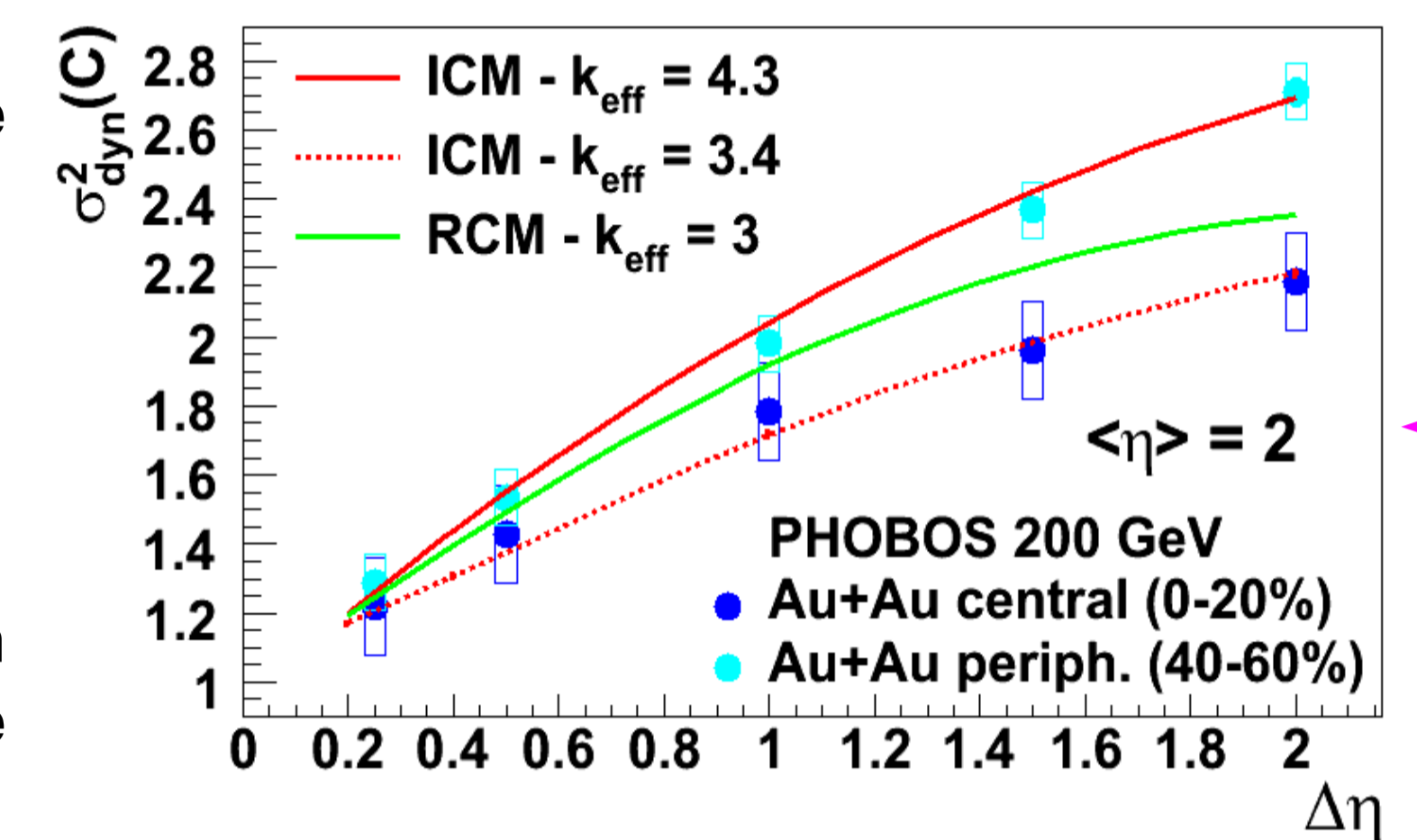
These values can be treated as the upper limits for k_{eff}

In **Resonance Cascade Model** for $\Delta(1900)$ we can change the number of charged final particles, but we show results only for $k_{eff} = 3$. For central collisions k_{eff} should be slightly smaller, for peripheral collisions - slightly larger

Using Forward-Backward fluctuations, the simultaneous determination of cluster width and size is difficult.

More direct measurement is possible in the studies of two-particle correlations.

Wei Li, "Two-particle angular correlations in pp, dA and AA collisions at PHOBOS", QM2006, Parallel Session 3.2



Assumptions for simple cluster models:

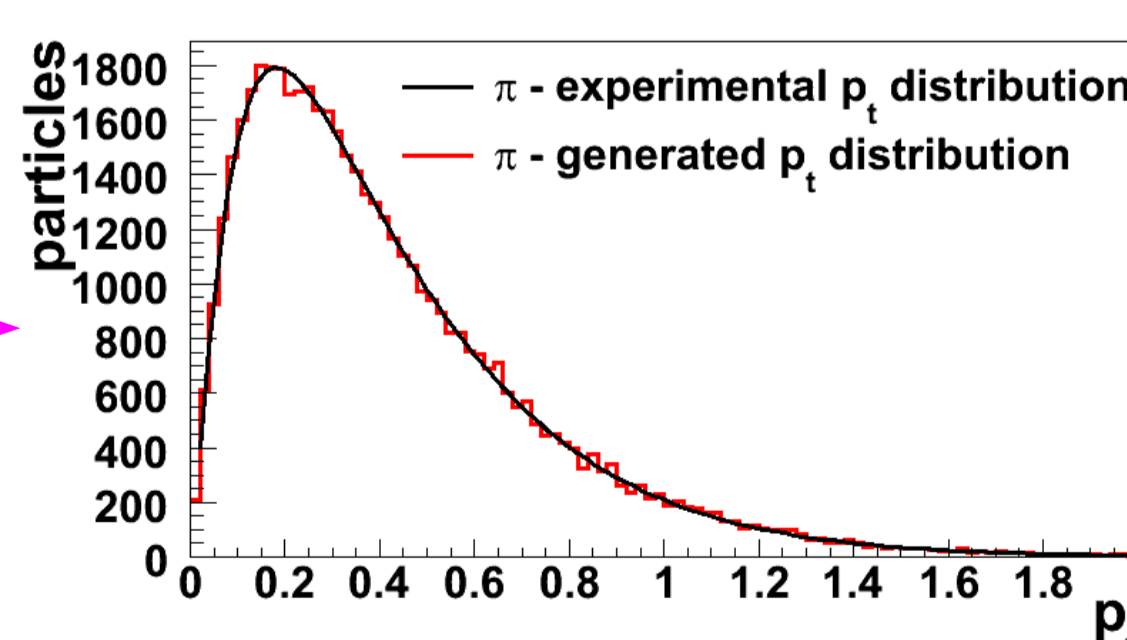
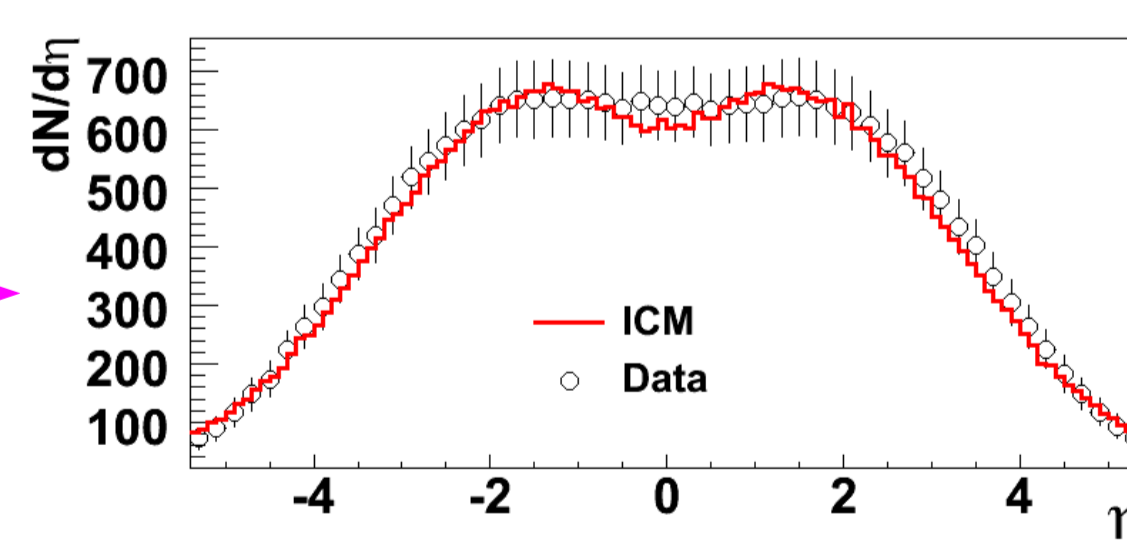
in Au+Au collisions intermediate objects (clusters) are produced and then decay into final particles
neutral particles (not measured in the experimental data compared) are neglected

Requirements:

- the transverse momentum distribution of produced particles is reproduced
- the pseudorapidity distribution of produced particles is reproduced

Isotropic Cluster Model:

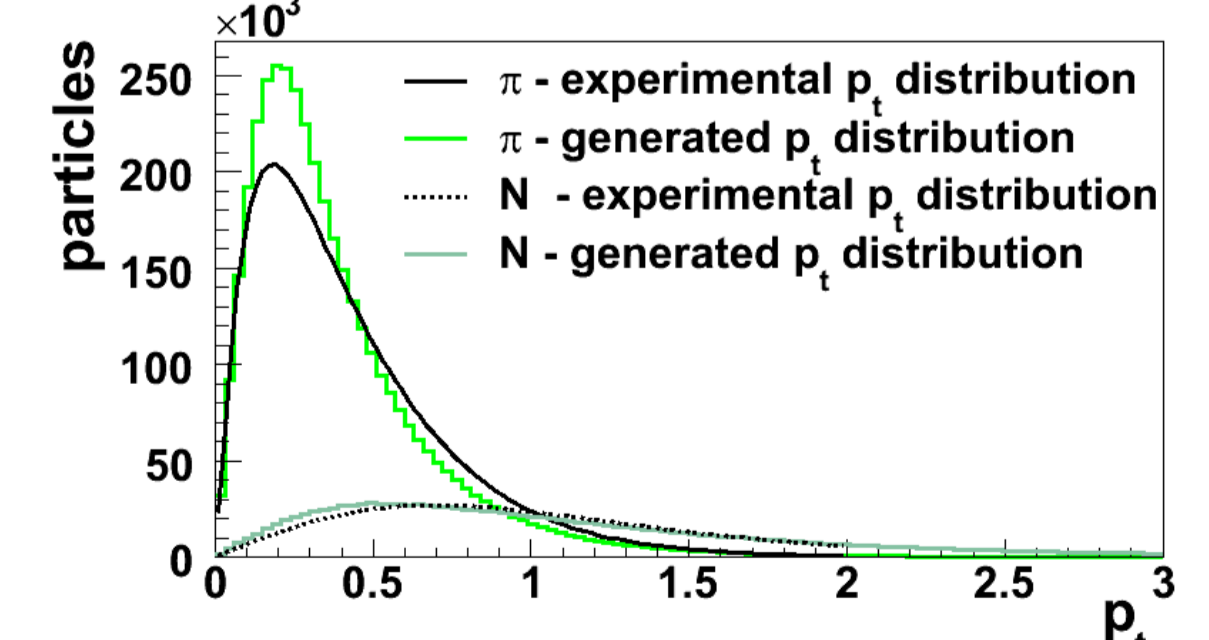
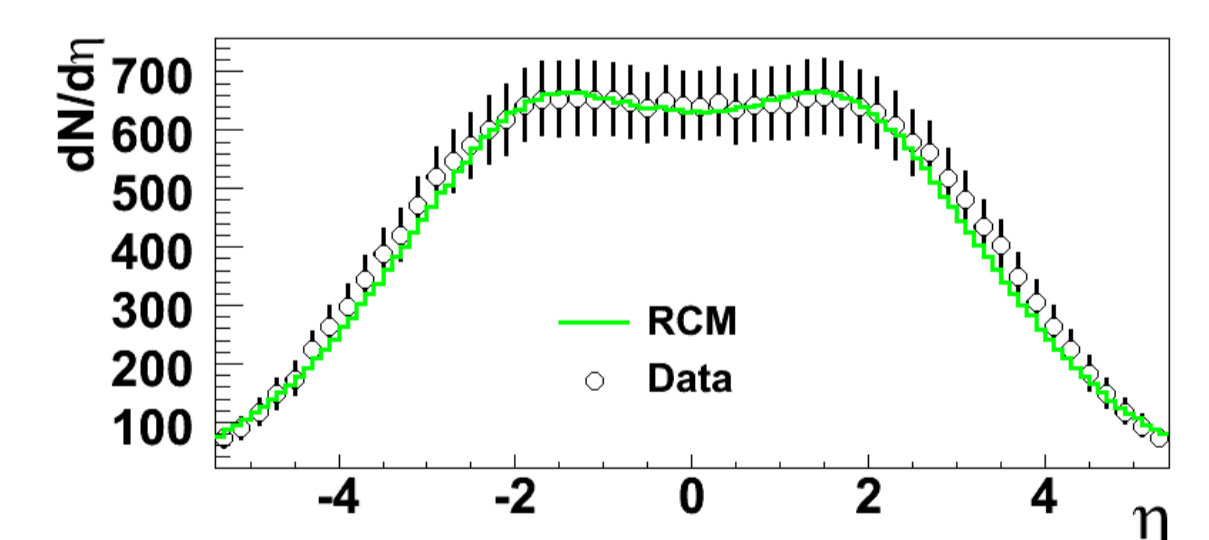
- clusters have negligibly small p_t
- clusters decay isotropically into pions
- p_t distribution of pions follows that from experiment
- p_z of initial cluster divided equally between decay products



Resonances Cascade Model:

- resonances may be a suitable source of clusters if they decay in several steps
- in decays $R(M_1) \rightarrow R(M_2) + \pi$, p_t of pion is determined by mass difference ($\langle p_t \rangle \approx 0.2$ GeV/c for $\Delta M \approx 0.3$ GeV)
- non zero initial p_t is necessary for reproduction of experimental p_t distribution

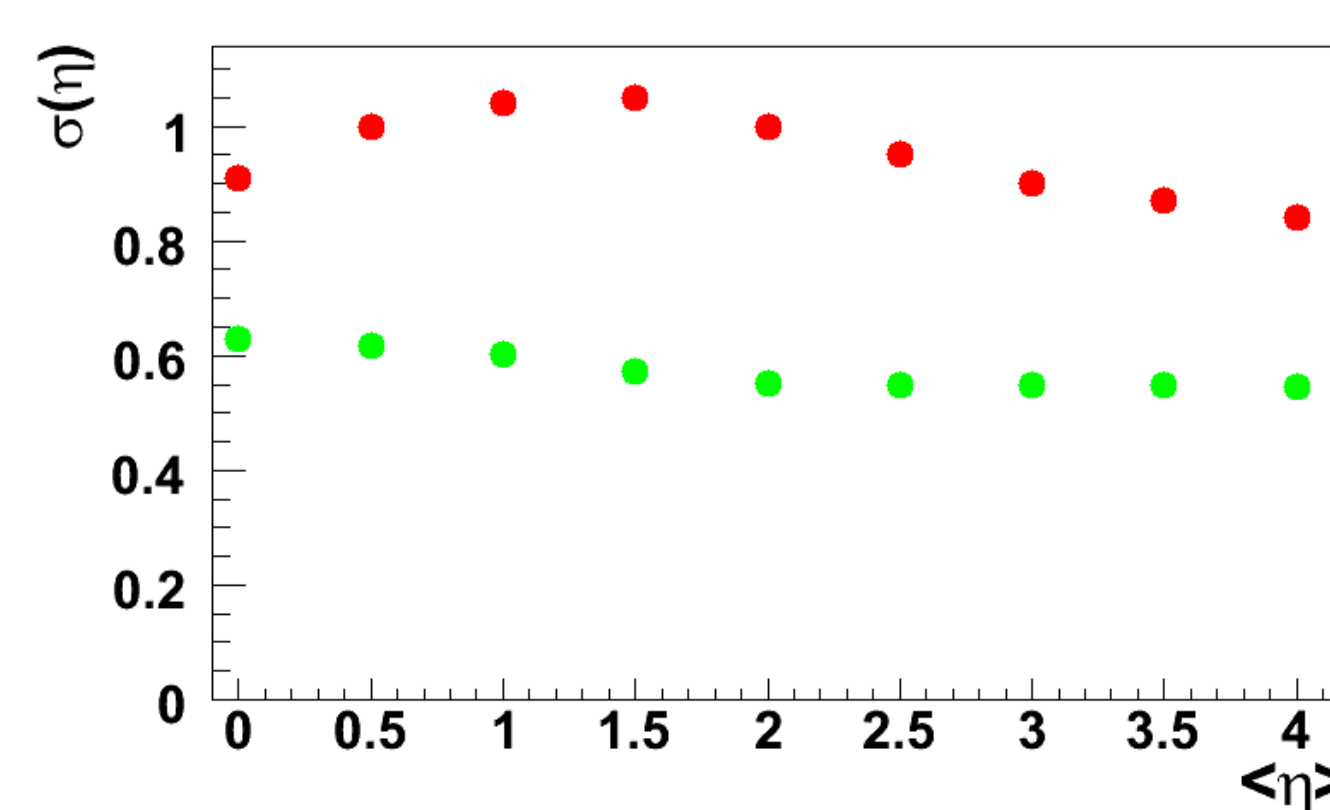
$\Delta(1900)$ with $\langle p_t \rangle = 2$ GeV/c is used as an example:
 $\Delta(1900) \rightarrow \Delta(1600) + \pi \rightarrow \Delta(1232) + 2\pi \rightarrow N + 3\pi$



Comparison of clusters width

Two cluster models present the extreme cases:

- in **Isotropic Cluster Model** the cluster width in pseudorapidity is the largest
- in **Resonance Cascade Model** for $\Delta(1900)$ the center of mass decay momenta of particles are small, initial p_t is large and thus cluster width is small



Studies of two-particle correlations give cluster width much smaller than 1, in favor of models with non-zero p_t of the clusters