PHOBOS collaboration



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A method for measuring **Elliptic Flow Fluctuations**

in 200 GeV Au+Au collisions at RHIC



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PHOBOS detector



Measuring elliptic flow fluctuations

200 GeV, Au+Au, track 200 GeV, Au+Au, hits 130 GeV, Au+Au, hits 62.4 GeV, Au+Au, hits 200 GeV, Cu+Cu, trac 200 GeV. Cu+Cu. hit 62.4 GeV, Cu+Cu, hit 22.4 GeV. Cu+Cu. h Au+Au Cu+Cu 300

Motivation



nucl:ex/0608025 Observed v₂ distribution

In this analysis v₂ is measured eventby-event: $g(v_2^{obs})$

The response of the event-by-event measurement, $K(v_2^{obs}, v_2)$, is calculated using detailed MC simulations, taking into account:



Au+Au, 200,130,62,4+19.6 GeV: PRL 94 122303 (2005) Cu+Cu, 200+62.4 GeV: nucl-ex/0610037 (sub.to PRL) Cu+Cu, 22.4 GeV: prel. QM06

Elliptic Flow in Cu+Cu collisions is observed to be significantly large.

Note: v_2 is quite large even for the most central collisions where the collision region is on average azimuthally symmetric relative to the impact parameter.

with respect to the minor axis of the ellipse defined by the collision zone.

Standard Eccentricity (ϵ_{std}) is calculated in a

Takes event-by-event fluctuations in the Glauber MC into account.

Participant Eccentricity (ϵ_{part}) is calculated

Ideal hydrodynamics predicts scaling between v_2 and ε .

The azimuthal anisotropy of the initial

Glauber MC with respect to the impact

Assumes no fluctuation in initial shape

collision region is quantified by the

eccentricity

parameter.

geometry for given b.

v_2 in Cu+Cu and Au+Au scale with ε_{nart} .

If $\varepsilon_{\text{part}}$ is the correct description of initial state geometry and if hydrodynamics work event by event, we should observe v_2 fluctuations

Finite-number fluctuations

Detector effects

Dependence of the resolution on multiplicity

The true v_2 distribution in data, $f(v_2)$, is calculated by finding a solution to:

$g(v_2^{obs}) = \int_0^\infty K(v_2^{obs}, v_2) f(v_2) dv_2$







The known η dependence of v₂ is used as input to the PDF

The likelihood of v_2^{obs} and ϕ_0 is defined as: $L(v_2^{obs}, \phi_0) = \prod_{i=1}^{n} P(\eta_i, \phi_i; v_2^{obs}, \phi_0)$

Event-by-event, v_2^{obs} and ϕ_0 are selected as the most likely values to generate the observed hit distribution.

Hit Distribution

Ф

The figure on the left shows the distribution of hits on the PHOBOS multiplicity array.

The array has a wide pseudorapitidy coverage:

-5.4 < η < 5.4

The detector has holes and granularity differences at midrapidity.

A new event-by-event flow measurement technique has been developed

to use all the available information in the detector and to allow an efficient correction for the acceptance effects.



$$\eta, \phi; v_2^{\text{obs}}, \phi_0) = \frac{1}{s(v_2^{\text{obs}}, \phi_0, \eta)} [1 + 2v_2(\eta) \cos(2\phi - 2\phi_0)]$$

where the normalization parameter, s, is included to make sure the PDF, folded by the acceptance, is normalized to the same value for different values of v_2^{obs} and ϕ_0 .

$$s(v_2^{obs}, \phi_0; \eta) = \int A(\eta, \phi) [1 + 2v_2(\eta) cos(2\phi - 2\phi_0)]$$

$$g(v_2^{obs}) = \int_0^\infty K(v_2^{obs}, v_2) f(v_2) dv_2$$

Determining the kernel $K(v_2^{obs}, v_2)$

P(

Extracting dynamical fluctuations $f(v_2)$

Event-by-event measurement $g(v_2^{obs})$

Probability dis 2 0 -4 -2 0 2 4 η

c flow ngle
$$\phi_0$$
 is n ϕ_0 ϕ_0

The probability to observe a hit at a certain

 $K(v_2^{obs}, v_2, n) = -$

fit

0.12

چ^{0.08} ک0.06

0.04

0.08

0.06

0.04

0.02 Au+Au, 200 GeV, h

PHOBOS preliminary

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The resolution of the measurement depends on the number of hits observed in the event.

Define n=observed number of hits

Therefore the kernel is $K(v_2^{obs}, v_2, n)$

The kernel is determined by "measuring" v_2^{obs} in bins of true v_2 and n in MC simulations.

HIJING, modified to include flow, is used to generate MC events. Particle azimuthal angles are redistributed with a probability distribution function defined by desired $v_2(\eta)$.

GEANT is used to simulate the detector response.

Plots on the right show $\langle v_2^{obs} \rangle$ and $\sigma(v_2^{obs})$ as a function of v_2 and n.

AMPT

Robustness of the kernel can be tested using events from a different MC generator, e.g. AMPT as "data" and reconstructing event-by-event flow fluctuations using kernel from HIJING.

 $v_2(\eta)$ and dN/d η were measured for AMPT and modified HIJING at the MC particle level.

AMPT is significantly different from HIJING.



Fitting $K(v_2^{obs}, v_2, n)$ with smooth functions reduces bin-to-bin fluctuations.

Theoretical distribution of $K(v_2^{obs}, v_2, n)$ modified for experimental effects is used as fit function

Assuming that the true v_2 distribution for a set of events in a given centrality class is independent of n, it is possible to integrate out the multiplicity $v_2 \rightarrow (An+B)v_2, \sigma = \frac{C}{\sqrt{2}} + D$ dependence: $K(v_2^{obs}, v_2) = \int K(v_2^{obs}, v_2, n) N(n) dn$



Results



In this measurement the quantities of interest are $\langle v_2 \rangle$ and $\sigma(v_2)$



Midrapidity <v₂> results from this event-by-event analysis show very good agreement with results from other event averaged PHOBOS analyses. Taking the ratio $\sigma(v_2)/\langle v_2 \rangle$ allows cancellation of various systematic errors.

These results are shown in the upper right plot in comparison with participant eccentricity predictions.

Results strongly support the existence of fluctuations in the initial state geometry and the event-by-event realization of the relationship $v_2 \propto \epsilon$.

See Talk by Constantin Loizides (Saturday, 4:20, Parallel 2.4)